

RESEARCH ARTICLE

10.1002/2013JA019396

Key Points:

- We observe the transmission of ULF waves from upstream to downstream
- The waves have the similar characteristics in both the upstream and downstream
- They have a large compressibility we conclude that these are magnetosonic waves

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Citation:

Shan, L., Q. Lu, M. Wu, X. Gao, C. Huang, T. Zhang, and S. Wang (2014), Transmission of large-amplitude ULF waves through a quasi-parallel shock at Venus, *J. Geophys. Res. Space Physics*, 119, 237–245, doi:10.1002/2013JA019396.

Received 1 SEP 2013

Accepted 22 DEC 2013

Accepted article online 27 DEC 2013

Published online 17 JAN 2014

Transmission of large-amplitude ULF waves through a quasi-parallel shock at Venus

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Abstract There exist large-amplitude ultralow frequency (ULF) waves in the upstream region of a quasi-parallel shock, which are excited due to the reflected ions by the shock. These waves are then brought back to the shock by the solar wind, and at last they coalesce and merge with the shock. In this paper, with the magnetic field measurements from Venus Express, for the first time we observe the transmission of large-amplitude ULF waves from the upstream region to the downstream under quasi-parallel shock conditions. These waves exist in both the upstream and downstream regions of the Venusian bow shock, which have the similar characteristics: their peak frequencies are 0.04–0.05 Hz in the spacecraft frame, their propagation angles do not change greatly, they have left-hand polarization with respect to the mean magnetic field in the spacecraft frame, and they also have a large compressibility. We conclude that they are magnetosonic waves. The generation mechanism of such waves at the Venusian bow shock is also discussed in the paper.

1. Introduction

A planetary bow shock is formed due to the interaction between the high-speed solar wind and the magnetosphere/ionosphere of a planet. A fast magnetosonic bow shock and its foreshock are significant features of a planet to the high-speed solar wind flow. The shock is an irreversible energy dissipation process, during which waves play an important role [Leroy *et al.*, 1982; Skopke *et al.*, 1983; Burgess *et al.*, 1989; Wilson *et al.*, 2007; Su *et al.*, 2012].

Due to the lack of an intrinsic planetary magnetic field, the interactions between the solar wind and the Venusian ionosphere have some special characteristics [Lu *et al.*, 2013]. Although the radius of Venus is similar to the Earth, the size of the Venusian bow shock is less than 1/10 of the Earth's bow shock [Slavin *et al.*, 1979]. The observations of Pioneer Venus Orbiter showed that the altitude of the subsolar bow shock is about $0.3 R_V$ ($1 R_V = 6051$ km, is the radius of Venus) [Zhang *et al.*, 1990] while it is about $14 R_E$ at the Earth. At the same time, in consideration of the small distance in dimension between the Venusian bow shock and the ionopause or induced magnetosphere, the neutral particles are easily leaked from the ionosphere and then ionized [Luhmann *et al.*, 1987]. Therefore, both the foreshock and magnetosheath of the Venusian bow shock have different characteristics from those of the Earth.

Numerous studies have revealed an abundance of wave activities or structures in the vicinity of the Venusian bow shock. Hoppe and Russell [1981] observed 1 Hz wave packets in the upstream region of the Venusian bow shock, and such waves were also identified in the upstream regions of the bow shocks at other solar system bodies [Fairfield, 1974; Orłowski *et al.*, 1990, 1992; Brain *et al.*, 2002; Le *et al.*, 2013]. Another wave phenomena detected at the upstream region of the Venusian bow shock is proton cyclotron waves [Delva *et al.*, 2008]. The waves propagate nearly (anti)parallel to the background magnetic field at a frequency just below the local proton cyclotron frequency, which are considered to be excited by the newly ionized particles escaped from the Venusian exosphere. A ring beam distribution is assumed to be formed for these pickup ions, and proton cyclotron waves are excited by the free energy associated with such a distribution [Wu and Davidson, 1972; Lu and Wang, 2006]. Proton cyclotron waves have also been observed at the upstream region of the Martian bow shock [Russell *et al.*, 1990; Brain *et al.*, 2002; Mazelle *et al.*, 2004; Wei *et al.*, 2011], and the ring beam distribution of pickup protons in the extended hydrogen corona of Mars has been verified by Barabash *et al.* [1991] with Phobos observations. On the basis of observations from Venus Express (VEX), Collinson *et al.* [2012] reported short large-amplitude magnetic structures (SLAMS) in the foreshock of the Venusian bow

shock, which are $\sim 1.5\text{--}11$ s in duration and have magnetic compression ratios between ~ 3 and 6. SLAMS were believed to steepen out of the large-amplitude ultralow frequency (ULF) waves in the upstream region of a quasi-parallel shock after they interact with diffuse ions, which had been observed at Earth, Jupiter, and comet Giacobini-Zinner [Schwartz *et al.*, 1992; Tsurutani *et al.*, 1990, 1993]. The ULF waves are suggested to be commonly associated with backstreaming ion distributions in the foreshock of a quasi-parallel shock [Paschmann *et al.*, 1979; Hoppe *et al.*, 1981; Thomsen *et al.*, 1985]. They attempt to propagate upstream; however, they are convected back toward the bow shock by the high-speed solar wind. Then, SLAMS grow out of these ULF waves when they are approaching the shock [Schwartz *et al.*, 1992; Scholer and Burgess, 1992; Tsubouchi and Lembège, 2004; Wilson *et al.*, 2009; Su *et al.*, 2012]. The nonlinear quadratic and cubic processes were found to play an important role during the evolution of SLAMS [Coca *et al.*, 2001; Zhu *et al.*, 2008].

There are also plenty of waves in the downstream region of the Venusian bow shock. Proton cyclotron waves [Russell *et al.*, 2006; Du *et al.*, 2009, 2010], as well as mirror-mode-like structures [Volwerk *et al.*, 2008], have been observed in the Venusian magnetosheath under quasi-perpendicular bow shock conditions. In the downstream region of a quasi-parallel shock, the magnetic fluctuations in the magnetosheath become more intense [Luhmann *et al.*, 1986; Du *et al.*, 2009, 2010], and the most likely source of these waves comes from the shock itself. Another important magnetic structure in the Venusian magnetosheath is the vortices observed by Venus Express (VEX), which are generated by the Kelvin-Helmholtz instability [Pope *et al.*, 2009]. Due to the lack of an intrinsic magnetic field at Venus, there is a direct contact between the fast-flowing solar wind and the Venusian ionosphere, and the velocity shear between the solar wind and the ions of ionopause is easy to be formed. The Kelvin-Helmholtz instability then grows in the Venusian magnetosheath [Pope *et al.*, 2009; Walker *et al.*, 2011], especially near the terminator region [Biernat *et al.*, 2007]. VEX has also observed coherent oscillations behind a quasi-perpendicular shock under the conditions with low Mach number and low β [Balikhin *et al.*, 2008], which are caused by the spatial pressure variations due to gyrations of the directly transmitted ions downstream of the ramp [Balikhin *et al.*, 2008; Ofman *et al.*, 2009].

In these previous studies, the wave characteristics at the Venusian bow shock are analyzed at either the upstream or the downstream region. Although Luhmann *et al.* [1986] indicated that the intense magnetic fluctuations in the downstream region of a quasi-parallel shock came from the shock itself, as we have known, the relations of the waves in the downstream and upstream regions have not been directly analyzed. In this paper, based on the measurement of the magnetic field from VEX, for the first time we reveal the existence of large-amplitude ULF waves at a quasi-parallel shock, and these waves can transmit the shock directly while keeping the wave characteristics, such as the frequency, propagating angle, and polarization. We first present two case studies, indicating that the large-amplitude ULF waves can transmit a quasi-parallel shock, and then the generation mechanism of such kind of waves is discussed.

2. Observations

VEX has an elliptical polar orbit with the periaapsis 250–300 km altitude. During a 24 h period orbit, the spacecraft encounters the inbound and outbound shock crossings [Barabash *et al.*, 2007]. The orbit covers large range of solar zenith angle of the Venusian bow shocks, which make it possible for the studying of subsolar bow shocks from valid measurements. The Venus Express magnetometer (VEX MAG) consists of two fluxgate sensors for a separation of magnetic effects of the spacecraft origin from the ambient space magnetic field [Zhang *et al.*, 2006].

2.1. Observations on 8 April 2011

The 15 min magnetic field data accompanied with large-amplitude ULF waves are obtained from VEX on 8 April 2011. In Figure 1, the measurements sampled at 1 Hz are displayed in the Venus Solar Orbital (VSO) coordinates. During the crossing from the upstream to the downstream of the Venusian bow shock, the average magnetic field changes from ~ 11 nT to ~ 25 nT. It is clear that VEX begins to encounter the shock at $\sim 03:11:00$ UT around $(1.31, -0.14, 0.46) R_V$ which is near the subsolar bow shock. Large-amplitude ULF waves are obviously observed in both the upstream and downstream regions with the peak-to-peak $\delta B/B \sim 0.3\text{--}3.1$. According to minimum variance analysis (MVA) method [Sonnerup and Cahill, 1967], the shock angle θ_{Bn} (the angle between the interplanetary magnetic field and the shock normal) can be calculated with the magnetic field during the interval from 03:10:00 UT to 03:13:00 UT, which is 29° , while it is 18° and 37° when the shock angle is calculated with magnetic coplanarity and the Venusian bow shock model [Zhang *et al.*, 2008; Shan

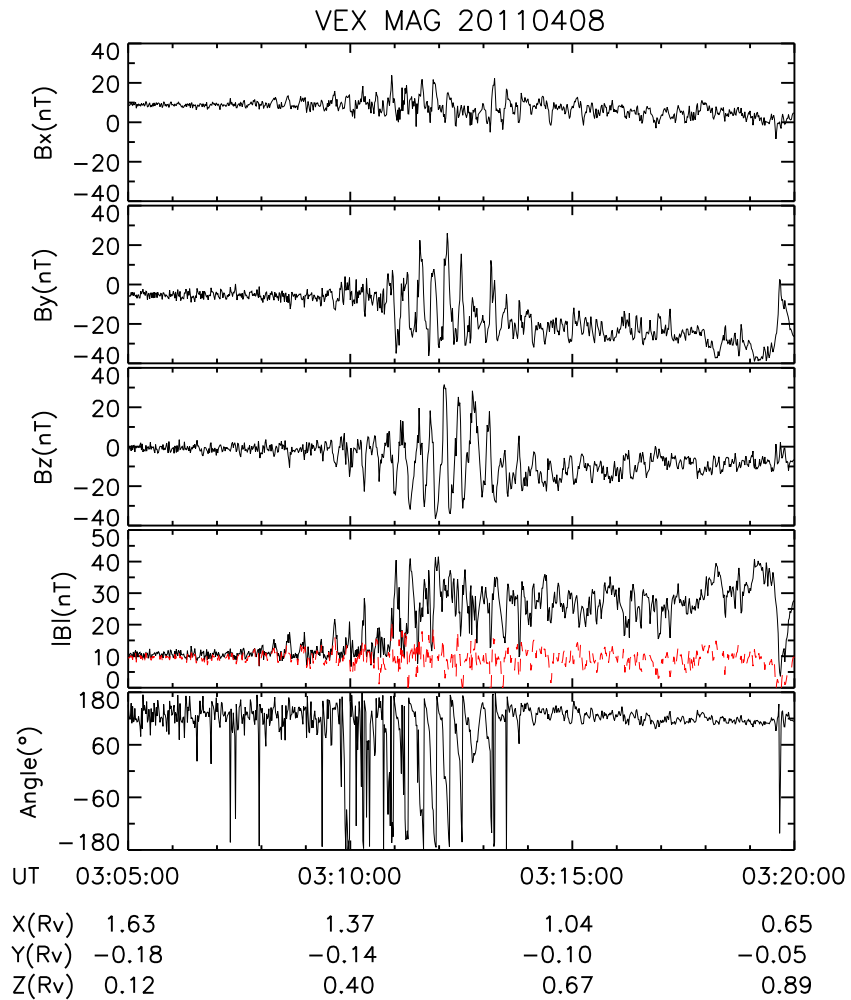


Figure 1. Three components of the magnetic field in VSO coordinates on 8 April 2011. The magnetic field data have 1 Hz time resolution. In the fourth panel, the solid line represents the total magnetic field in VSO coordinates, while the red dashed line is the projection of the magnetic field on the shock normal direction determined with the MVA method. The azimuthal angle of the magnetic field in the plane perpendicular to the minimum variance direction is shown in the fifth panel.

et al., 2013], respectively. Obviously, it is a quasi-parallel shock. A low-pass filter is utilized before we calculate the shock normal. In Figure 1 (fourth panel), we also plot the projection of the magnetic field on the shock normal determined with the MVA method, and values in the upstream are approximately equal to those in the downstream, which validates the identified shock normal. In Figure 1 (fifth panel), we show the azimuthal angle of the magnetic field in the plane perpendicular to the minimum variance direction. Here in order to reduce the disturbance from the shock structure, we select the downstream period from 03:11:50 to 03:13:20 UT to determine the variance reference frame. Then, the data between 03:05:00 and 03:20:00 is transferred into the variance reference frame, and the azimuthal angle is calculated. The ratios of the maximum to intermediate and intermediate to minimum eigenvalues are 1.8 and 15.5, respectively, and it demonstrates that the analysis process is reliable. In Figure 1 (fifth panel), the angle changes from 180° to -180° monotonously, and then a jump is followed, which indicates the elliptical polarization of the waves. It is obvious that the amplitude decreases as the waves propagate to the farther distance from the shock in the magnetosheath. Moreover, the waves disappear at ~03:15:00 UT while the satellite reaches the ionopause at ~03:19:30 UT.

Figure 2 shows the power spectrum of the ULF waves in the upstream (03:06:30–03:10:40 UT) and downstream (03:12:20–03:16:30 UT) regions on 8 April 2011, separated by the transverse (red) and compressional (black) power, respectively. In Figure 2, the blue dashed lines represent the local proton cyclotron frequency $f_{cp} = qB / (2\pi m_p)$, where B , q , and m_p are the magnetic field, charge, and mass of the proton. The peak frequency of the

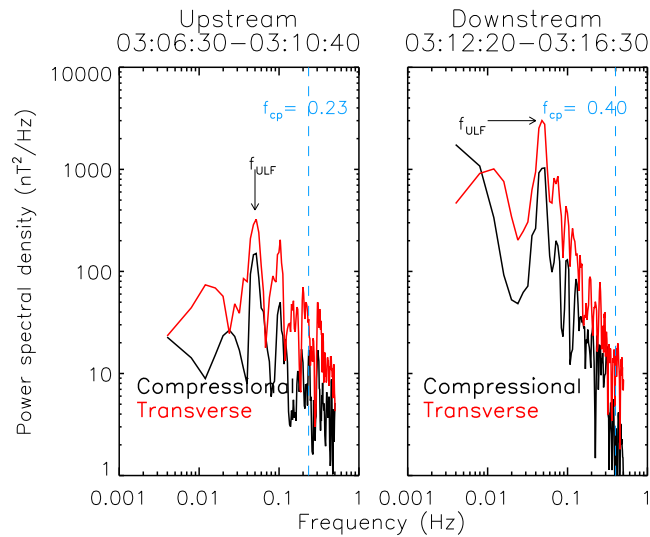


Figure 2. The power spectrum of the ULF waves upstream and downstream of the Venusian bow shock on 8 April 2011, separated in transverse and compressional power, respectively.

ULF waves (the waves have the maximum amplitude at this frequency) is f_{ULF} . Obviously, the transverse component dominates the power spectrum. The peak frequency of the upstream waves is almost the same as that of the downstream waves, which is about 0.05 Hz. In order to investigate the quasi-monochromatic ULF waves in more detail, the principal axis analysis [Rankin and Kurtz, 1970; McPherron et al., 1972] is used to determine the characteristics of the ULF waves upstream and downstream of the shock. Figure 3 displays the hodograms of the wave magnetic field in the principal axis coordinates (B_i , B_j , and B_k) during the intervals 03:09:50–03:10:40 (upstream) and 03:13:00–03:13:50 UT (downstream) on 8 April 2011, respectively, where \mathbf{k} is the computed propagation direction of the waves, and it may be parallel or antiparallel to the wave vector. In the left boxes, the mean magnetic field points out to the plane; therefore, the waves exhibit left-hand elliptical polarization with respect to the mean magnetic field in the spacecraft frame. The propagation angle θ_{kB} (the angle between \mathbf{k} and the mean magnetic field) and ellipticity ε can be calculated with the quadrature power spectral matrix analysis method [Means, 1972; Song and Russell, 1999]. In the upstream and the downstream regions of the

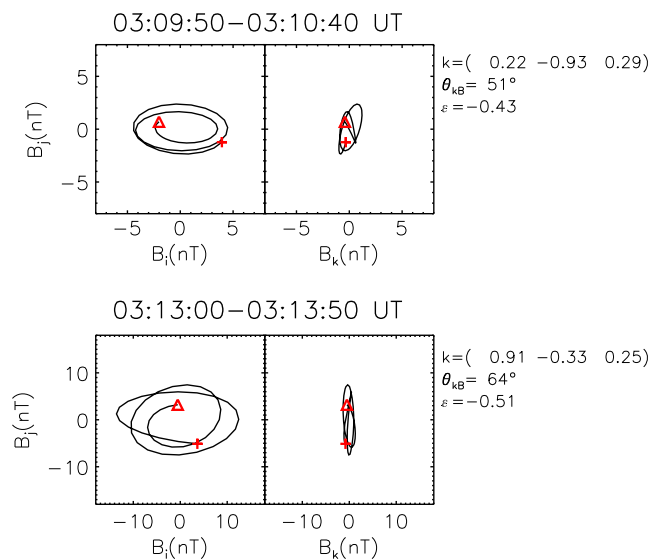


Figure 3. The hodograms of the wave magnetic field in the principal axis coordinates (B_i , B_j , and B_k) during the intervals 03:09:50–03:10:40 (upstream) and 03:13:00–03:13:50 (downstream) on 8 April 2011, respectively, where \mathbf{k} is the computed propagation direction of the waves and it may be parallel or antiparallel to the wave vector. In the figure, “plus” and “triangle” are the beginning and end of the interval.

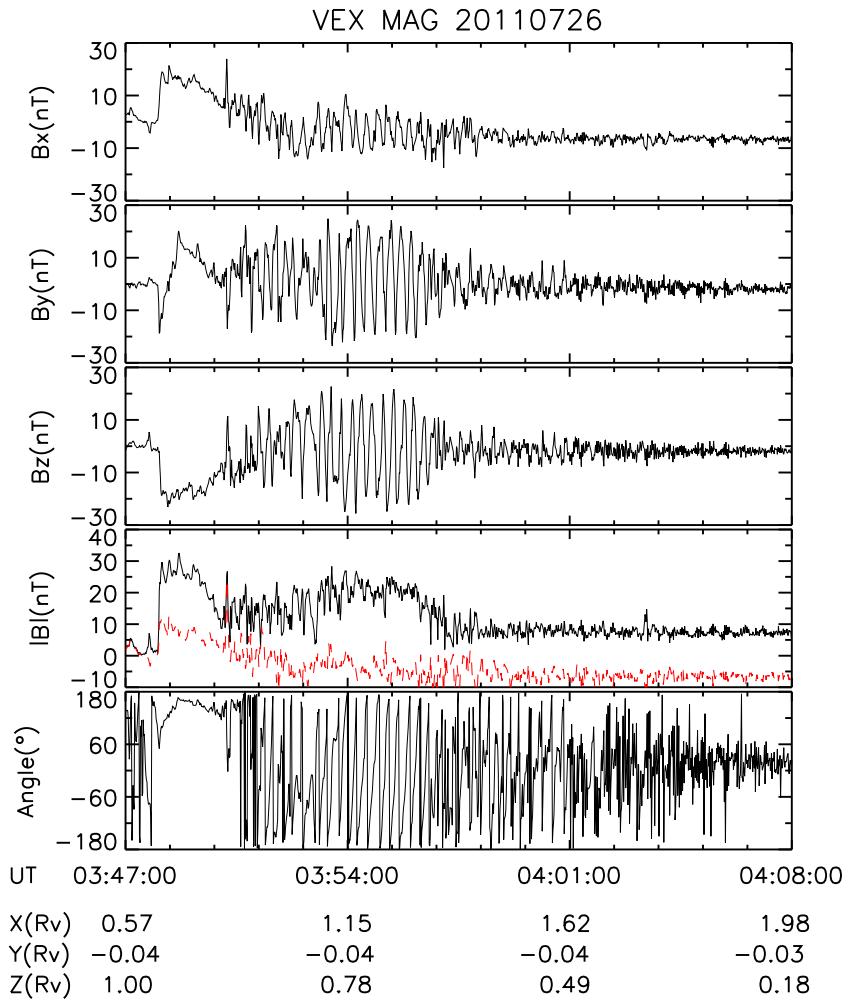


Figure 4. Same as Figure 1 but for observations on 26 July 2011.

shock, θ_{kB} is about 51° and 64° , while ε is -0.43 and -0.51 , respectively. The minus sign of ε means that the waves are left-hand polarized with respect to the mean magnetic field, while the plus sign is for right-hand polarization. Therefore, the waves exhibit left-hand polarization with respect to the mean magnetic field in the spacecraft frame. However, these waves are propagating in the high-speed solar wind, and they may have right-hand polarization in the plasma frame due to the Doppler shift [Hoppe and Russell, 1983; Mazelle et al., 2003].

2.2. Observations on 26 July 2011

The similar large-amplitude waves are also identified on 26 July 2011. The 21 min magnetic field components and the azimuthal angle are displayed in Figure 4. Obviously, VEX is on the outbound crossing of its trajectory, and it detects the ionopause at $\sim 03:48:00$ UT and enters the upstream region of the shock at $\sim 03:57:20$ UT around $(1.38, -0.04, 0.66) R_V$. The shock angle θ_{Bn} is 10° , 28° , and 4° determined by MVA (based on the interval from $03:55:00$ UT to $03:58:00$ UT), magnetic coplanarity, and the Venusian bow shock model, respectively, suggesting that it is a quasi-parallel shock. It is noted that the large-amplitude waves are observed in the upstream region between $03:57:20$ and $04:08:00$ UT, corresponding to the position from $(1.38, -0.04, 0.66)$ to $(1.98, -0.03, 0.18) R_V$. The ULF waves have peak-to-peak $\delta B/B \sim 0.2-2.7$. The downstream measurements during the period of $03:53:00-03:55:00$ UT are selected to determine the azimuthal angle of the magnetic field in the plane perpendicular to the minimum variance direction. The analysis is valid because the ratios of the maximum to intermediate and intermediate to minimum eigenvalues are 1.1 and 30.5, respectively. For the azimuthal angle, we can obviously observe a monotonous change from 180° to -180° , and then a jump is followed, which indicates the elliptical polarization of the waves.

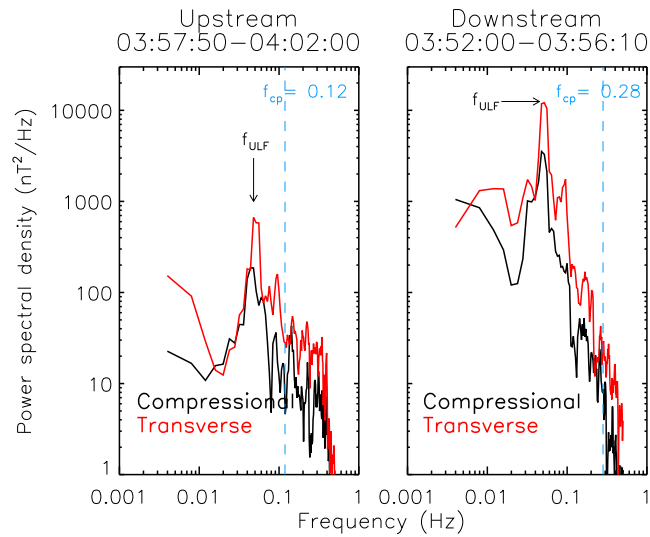


Figure 5. Same as Figure 2 but for observations on 26 July 2011.

Figure 5 shows the power spectrum of the large-amplitude waves on 26 July 2011. The peak frequency of the upstream waves, ~ 0.05 Hz in the spacecraft frame, is similar to that observed on the downstream side of the shock. Therefore, it is shown that the quasi-monochromatic waves are possible generated at the upstream region of the shock and then convected into downstream by high-speed solar wind. The hodograms of the wave magnetic field upstream (04:00:10–04:01:00 UT) and downstream (03:55:20–03:56:10 UT) of the shock are displayed in Figure 6. The average magnetic field points out to the plane in the left boxes. Hence, the waves exhibit left-hand elliptical polarization with respect to the mean magnetic field in the spacecraft frame. The propagation angle θ_{kB} varies from 31° (upstream) to 41° (downstream), and the ellipticity ε is -0.55 and -0.90 , respectively.

3. Conclusions and Discussion

This paper presents the observations of the transmission of large-amplitude ULF waves from the upstream region into the downstream region under the quasi-parallel shock conditions in the vicinity of the Venusian bow shock. Two events, on 8 April 2011 and 26 July 2011, are analyzed in detail. By examining the magnetic field data

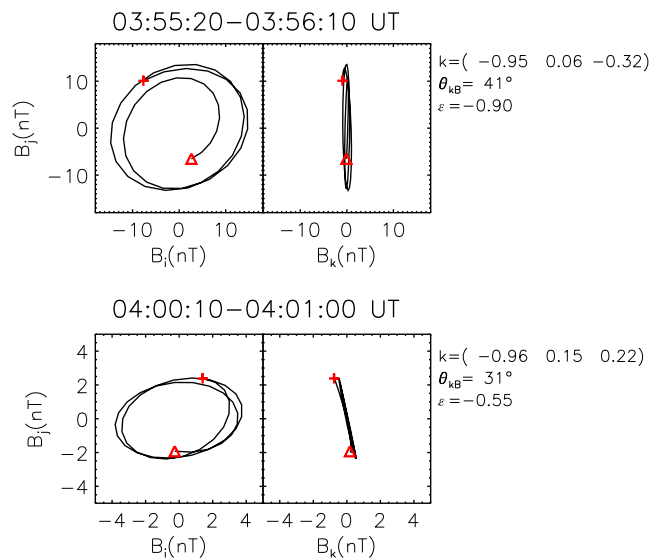


Figure 6. Same as Figure 3 but for observations on 26 July 2011.

Table 1. The Characteristics of the ULF Waves

Time (UT)	θ_{Bn} (deg)	θ_{kB} (deg)		Ellipticity ε		Peak Frequency (Hz)		$\delta B/B$
		Upstream	Downstream	Upstream	Downstream	Upstream	Downstream	
22 Sep 2006 041230–042230	44	42	59	−0.61	−0.46	0.044	0.040	0.5–3.2
27 Nov 2008 090730–091730	18	15	49	−0.75	−0.88	0.042	0.042	0.5–2.4
8 Apr 2011 030630–031630	29	51	64	−0.43	−0.51	0.050	0.050	0.3–3.1
19 May 2011 020000–021000	32	42	41	−0.60	−0.76	0.043	0.047	0.3–2.4
26 July 2011 035200–040200	10	31	41	−0.55	−0.90	0.050	0.054	0.2–2.7

measured by VEX, we find five similar cases, and their characteristics are listed in Table 1. In these cases, the large-amplitude quasi-monochromatic waves exist in both the upstream and the downstream regions of the Venusian bow shock. The characteristics of the downstream waves are similar to that of the upstream waves, except with a larger amplitude, and these characteristics include the following: left-hand polarization with respect to the mean magnetic field in the spacecraft frame, the peak frequency 0.04–0.05 Hz in the spacecraft frame, and the stable propagating angle (except the event on 27 November 2007, the propagating angle changed from 15° in the upstream to 49° in the downstream). We conclude that these waves are magnetosonic waves.

We can rule the vortices and coherent oscillations out of the generation mechanism for these waves. The vortices is formed by Kelvin-Helmholtz instability due to the velocity shear between solar wind and ionospheric ions [Pope *et al.*, 2009; Walker *et al.*, 2011], and they only exist in the downstream region of the Venusian bow shock. The coherent oscillations are caused by the spatial pressure variations due to gyrations of the directly transmitted ions downstream of the ramp [Balikhin *et al.*, 2008; Ofman *et al.*, 2009], and they only exist in the downstream region of a quasi-perpendicular shock with low Mach number and low β [Balikhin *et al.*, 2008; Ofman *et al.*, 2009]. These waves observed in this paper may be excited by the reflected ions by a quasi-parallel shock through electromagnetic ion beam instabilities [Winske and Gary, 1986; Gary, 1991], and they are then convected toward the shock by the high-speed solar wind. However, these waves will steepen into SLAMS after they interact with diffuse ions. In the paper, we observe large-amplitude ULF waves in both the upstream and downstream regions of the Venusian bow shock under quasi-parallel shock conditions, and they have the similar characteristics. These waves are considered to be able to transmit from the upstream region to the downstream. One possibility is that the size of the Venusian foreshock is much smaller than that of Earth, and the region with diffuse ions is also small. Therefore, the waves do not have enough time to grow into SLAMS when they propagate through the Venusian foreshock.

Luhmann *et al.* [1983] observed large-amplitude magnetic fluctuations with a period of 10–40 s in the downstream region of a quasi-parallel shock and indicated that the source of these waves was from the upstream region of the shock. Our study shows that the large-amplitude ULF waves in the upstream region can transmit the Venusian bow shock under the quasi-parallel shock conditions, and it gives the direct evidence that the intense downstream waves behind a quasi-parallel shock can be from the upstream region.

Acknowledgments

The authors thank all members of the MAG team for substantial and accuracy measurements. We are grateful for discussions with Y. Lin, G. Le, X. Blanco-Cano, and G. A. Collinson. This research was supported by the 973 Program (2012CB825602 and 2013CBA01503), and the National Science Foundation of China, grants 41174124, 41274144, 41331067, 41121003, Ocean Public Welfare Scientific Research Project, State Oceanic Administration People's Republic of China (201005017), and CAS Key Research Program KZZD-EW-01. Philippa Browning thanks Jian Du and an anonymous reviewer for their assistance in evaluating this paper.

References

- Balikhin, M. A., T. L. Zhang, M. Gedalin, N. Y. Ganushkina, and S. A. Pope (2008), Venus Express observes a new type of shock with pure kinematic relaxation, *Geophys. Res. Lett.*, *35*, L01103, doi:10.1029/2007GL032495.
- Barabash, S., E. Dubinin, N. Pissarenko, R. Lundin, and C. T. Russell (1991), Picked-up protons near Mars: Phobos observations, *Geophys. Res. Lett.*, *18*, 1805–1808, doi:10.1029/91GL02082.
- Barabash, S., *et al.* (2007), The Analyser of Space Plasma and Energetic Atoms (ASPERA-4) for the Venus Express mission, *Planet. Space Sci.*, *55*, 1772–1792.
- Biernat, H. K., N. V. Erkaev, U. V. Amerstorfer, T. Penz, and H. I. M. Lichtenegger (2007), Solar wind flow past Venus and its implications for the occurrence of the Kelvin-Helmholtz instability, *Planet. Space Sci.*, *55*, 1793–1803, doi:10.1016/j.pss.2007.01.006.
- Brain, D. A., F. Bagenal, M. H. Acuna, J. E. P. Connerney, D. H. Crider, C. Mazelle, D. L. Mitchell, and N. F. Ness (2002), Observations of low frequency electromagnetic plasma waves upstream from the Martian shock, *J. Geophys. Res.*, *107*(A6), 1076, doi:10.1029/2000JA000416.
- Burgess, D., W. P. Wilkinson, and S. J. Schwartz (1989), Ion distributions and thermalization at perpendicular and quasi-perpendicular supercritical collisionless shocks, *J. Geophys. Res.*, *94*(A7), 8783–8792, doi:10.1029/JA094iA07p08783.
- Coca, D., M. A. Balikhin, S. A. Billings, H. S. C. Alleyne, and M. Dunlop (2001), Time domain analysis of plasma turbulence observed upstream of a quasi-parallel shock, *J. Geophys. Res.*, *106*(A11), 25,005–25,021, doi:10.1029/2000JA000431.

- Collinson, G. A., L. B. Wilson, D. G. Sibeck, N. Shane, T. L. Zhang, T. E. Moore, A. J. Coates, and S. Barabash (2012), Short large-amplitude magnetic structures (SLAMS) at Venus, *J. Geophys. Res.*, *117*, A10221, doi:10.1029/2012JA017838.
- Delva, M., T. L. Zhang, M. Volwerk, W. Magnes, C. T. Russell, and H. Y. Wei (2008), First upstream proton cyclotron wave observations at Venus, *Geophys. Res. Lett.*, *35*, L03105, doi:10.1029/2007GL032594.
- Du, J., T. L. Zhang, C. Wang, M. Volwerk, M. Delva, and W. Baumjohann (2009), Magnetosheath fluctuations at Venus for two extreme orientations of the interplanetary magnetic field, *Geophys. Res. Lett.*, *36*, L09102, doi:10.1029/2009GL037725.
- Du, J., T. L. Zhang, W. Baumjohann, C. Wang, M. Volwerk, Z. Vörös, and L. Guicking (2010), Statistical study of low-frequency magnetic field fluctuations near Venus under the different interplanetary magnetic field orientations, *J. Geophys. Res.*, *115*, A12251, doi:10.1029/2010JA015549.
- Fairfield, D. H. (1974), Whistler waves observed upstream from collisionless shocks, *J. Geophys. Res.*, *79*, 1368–1378, doi:10.1029/JA079i010p01368.
- Gary, S. P. (1991), Electromagnetic ion/ion instabilities and their consequences in space plasmas: A review, *Space Sci. Res.*, *56*, 373–415.
- Hoppe, M. M., and C. T. Russell (1981), On the nature of ULF waves upstream of planetary bow shocks, *Adv. Space Res.*, *1*, 327–332.
- Hoppe, M. M., and C. T. Russell (1983), Plasma rest frame frequencies and polarizations of the low-frequency upstream waves: ISEE 1 and 2 observations, *J. Geophys. Res.*, *88*, 2021–2028.
- Hoppe, M. M., C. T. Russell, L. A. Frank, T. E. Eastman, and E. W. Greenstadt (1981), Upstream hydromagnetic waves and their association with backstreaming ion populations: ISEE-1 and -2 observations, *J. Geophys. Res.*, *86*, 4471–4492.
- Le, G., P. J. Chi, X. Blanco-Cano, S. Boardsen, J. A. Slavin, and B. J. Anderson (2013), Upstream ultra-low frequency waves in Mercury's foreshock region: MESSENGER magnetic field observations, *J. Geophys. Res. Space Physics*, *118*, 2809–2823, doi:10.1002/jgra.50342.
- Leroy, M. M., D. Winske, C. C. Goodrich, C. S. Wu, and K. Papadopoulos (1982), The structure of perpendicular bow shocks, *J. Geophys. Res.*, *87*(A7), 5081–5094, doi:10.1029/JA087iA07p05081.
- Lu, Q. M., and S. Wang (2006), Electromagnetic waves downstream of quasi-perpendicular shocks, *J. Geophys. Res.*, *111*, A05204, doi:10.1029/2005JA011319.
- Lu, Q. M., L. C. Shan, T. L. Zhang, G. P. Zank, Z. W. Yang, M. Y. Wu, A. M. Du, and S. Wang (2013), The role of pickup ions on the structure of the Venusian bow shock and its implication for the termination shock, *Astrophys. J.*, *773*, L24, doi:10.1088/2041-8205/773/L24.
- Luhmann, J. G., M. Tatrallyay, C. T. Russell, and D. Winterhalter (1983), Magnetic field fluctuations in the Venus magnetosheath, *Geophys. Res. Lett.*, *10*, 655–658, doi:10.1029/GL010i008p0655.
- Luhmann, J. G., C. T. Russell, and R. C. Elphic (1986), Spatial distributions of magnetic field fluctuations in the dayside magnetosheath, *J. Geophys. Res.*, *91*(A2), 1711–1715, doi:10.1029/JA091iA02p01711.
- Luhmann, J. G., C. T. Russell, J. L. Phillips, and A. Barns (1987), On the role of the quasi-parallel bow shock in ion pickup: A lesson from Venus?, *J. Geophys. Res.*, *92*(A3), 2544–2550, doi:10.1029/JA092iA03p02544.
- Mazelle, C., et al. (2003), Production of gyrating ions from nonlinear wave-particle interaction upstream from the Earth's bow shock: A case study from Cluster-CIS, *Planet. Space Sci.*, *51*, 785–795.
- Mazelle, C., et al. (2004), Bow shock and upstream phenomena at Mars, *Space Sci. Res.*, *111*, 115–181.
- McPherron, R. L., C. T. Russell, and J. Coleman (1972), Fluctuating magnetic fields in the magnetosphere. II: ULF waves, *Space Sci. Rev.*, *13*(3), 411–454, doi:10.1007/BF00219165.
- Means, J. D. (1972), Use of the three-dimensional covariance matrix in analyzing the polarization properties of plane waves, *J. Geophys. Res.*, *77*, 5551–5559, doi:10.1029/JA077i028p05551.
- Ofman, L., M. Balikhin, C. T. Russell, and M. Gedalin (2009), Collision-less relaxation of ion distributions downstream of laminar quasi-perpendicular shocks, *J. Geophys. Res.*, *114*, A09106, doi:10.1029/2009JA014365.
- Orlowski, D. S., G. K. Crawford, and C. R. Russell (1990), Upstream waves at mercury, Venus and Earth: Comparison of the properties of one Hertz waves, *Geophys. Res. Lett.*, *17*(13), 2293–2296, doi:10.1029/GL017i013p02293.
- Orlowski, D. S., C. R. Russell, and R. P. Lepping (1992), Wave phenomena in the upstream region of Saturn, *J. Geophys. Res.*, *97*, 19,187–19,199, doi:10.1029/92JA01461.
- Paschmann, G., N. Sckopke, S. J. Bame, J. R. Asbridge, J. T. Gosling, C. T. Russell, and E. W. Greenstadt (1979), Association of low frequency waves with suprathermal ions in the upstream solar wind, *Geophys. Res. Lett.*, *6*, 209–212, doi:10.1029/GL006i003p0209.
- Pope, S. A., M. A. Balikhin, T. L. Zhang, A. O. Fedorov, M. Gedalin, and S. Barabash (2009), Giant vortices lead to ion escape from Venus and redistribution of plasma in the ionosphere, *Geophys. Res. Lett.*, *36*, L07202, doi:10.1029/2008GL036977.
- Rankin, D., and R. Kurtz (1970), Statistical study of micropulsation polarizations, *J. Geophys. Res.*, *75*, 5444–5458, doi:10.1029/JA075i028p05444.
- Russell, C. T., J. G. Luhmann, K. Schwingenschuh, W. Riedler, and Y. Yeroshenko (1990), Upstream waves at Mars: Phobos observations, *Geophys. Res. Lett.*, *17*, 897–900.
- Russell, C. T., S. S. Mayerberger, and X. Blanco-Cano (2006), Proton cyclotron waves at Mars and Venus, *Adv. Space Res.*, *38*, 745–751.
- Scholer, M., and D. Burgess (1992), The role of upstream waves in supercritical quasi-parallel shock re-formation, *J. Geophys. Res.*, *97*(A6), 8319–8326, doi:10.1029/92JA00312.
- Schwartz, S. J., D. Burgess, W. P. Wilkinson, R. L. Kessel, M. Dunlop, and H. Luehr (1992), Observations of short large-amplitude magnetic structures at a quasi-parallel shock, *J. Geophys. Res.*, *97*, 4209–4227.
- Sckopke, N., G. Paschmann, S. J. Bame, J. T. Gosling, and C. T. Russell (1983), Evolution of ion distributions across the nearly perpendicular bow shock: Specularly and non-specularly reflected-gyrating ions, *J. Geophys. Res.*, *88*(A8), 6121–6136, doi:10.1029/JA088iA08p06121.
- Shan, L. C., Q. M. Lu, T. L. Zhang, X. L. Gao, C. Huang, Y. Q. Su, and S. Wang (2013), Comparison between magnetic coplanarity and MVA methods in determining the normal of Venusian bow shock, *Chin. Sci. Bull.*, *58*, 2469–2472, doi:10.1007/s11434-013-5675-8.
- Slavin, J. A., R. C. Elphic, C. T. Russell, J. H. Wolfe, and D. S. Intriligator (1979), Position and shape of the Venus bow shock: Pioneer Venus Orbiter observations, *Geophys. Res. Lett.*, *6*, 901–904, doi:10.1029/GL006i011p00901.
- Song, P., and C. T. Russell (1999), Time series data analyses in space plasmas, *Space Sci. Res.*, *87*, 387–463.
- Sonnerup, B. U. Ö., and L. J. Cahill Jr. (1967), Magnetopause structure and attitude from explorer 12 observations, *J. Geophys. Res.*, *72*, 171–183, doi:10.1029/JZ072i001p0171.
- Su, Y., Q. Lu, C. Huang, M. Wu, X. Gao, and S. Wang (2012), Particle acceleration and generation of diffuse superthermal ions at a quasi-parallel collisionless shock: Hybrid simulations, *J. Geophys. Res.*, *117*, A08107, doi:10.1029/2012JA017736.
- Thomsen, M. F., J. T. Gosling, and S. J. Bame (1985), Gyrating ions and large-amplitude monochromatic MHD waves upstream of the Earth's bow shock, *J. Geophys. Res.*, *90*, 267–273.
- Tsubouchi, K., and B. Lembège (2004), Full particle simulations of short large-amplitude magnetic structures (SLAMS) in quasi-parallel shocks, *J. Geophys. Res.*, *109*, A02114, doi:10.1029/2003JA010014.

- Tsurutani, B. T., E. J. Smith, H. Matsumoto, A. L. Brinca, and N. Omid (1990), Highly nonlinear magnetic pulses at comet Giacobini-Zinner, *Geophys. Res. Lett.*, *17*, 757–760, doi:10.1029/GL017i006p00757.
- Tsurutani, B. T., D. J. Southwood, E. J. Smith, and A. Balogh (1993), A survey of low frequency waves at Jupiter: The Ulysses encounter, *J. Geophys. Res.*, *98*(A12), 21,203–21,216, doi:10.1029/93JA02586.
- Volwerk, M., T. L. Zhang, M. Delva, Z. Vörös, W. Baumjohann, and K.-H. Glassmeier (2008), Mirror-mode-like structures in Venus' induced magnetosphere, *J. Geophys. Res.*, *113*, E00B16, doi:10.1029/2008JE003154.
- Walker, S. N., M. A. Balikhin, T. L. Zhang, M. E. Gedalin, S. A. Pope, A. P. Dimmock, and A. O. Fedorov (2011), Unusual nonlinear waves in the Venusian magnetosheath, *J. Geophys. Res.*, *116*, A01215, doi:10.1029/2010JA015916.
- Wei, H. Y., C. T. Russell, T. L. Zhang, and X. Blanco-Cano (2011), Comparative study of ion cyclotron waves at Mars, Venus and Earth, *Planet. Space Sci.*, *59*, 1039–1047, doi:10.1016/j.pss.2010.01.004.
- Wilson, L. B., III, C. Cattell, P. J. Kellogg, K. Goetz, K. Kersten, L. Hanson, and R. MacGregor (2007), Waves in interplanetary shocks: A Wind/WAVES study, *Phys. Rev. Lett.*, *99*, 041101.
- Wilson, L. B., III, C. A. Cattell, P. J. Kellogg, K. Goetz, K. Kersten, J. C. Kasper, A. Szabo, and K. Meziane (2009), Low-frequency whistler waves and shocklets observed at quasi-perpendicular interplanetary shocks, *J. Geophys. Res.*, *114*, A10106, doi:10.1029/2009JA014376.
- Winske, D., and S. P. Gary (1986), Electromagnetic beam instabilities driven by cool heavy ion beams, *J. Geophys. Res.*, *91*(A6), 6825–6832, doi:10.1029/JA091iA06p06825.
- Wu, C. S., and R. C. Davidson (1972), Electromagnetic instabilities produced by neutral-particle ionization in interplanetary space, *J. Geophys. Res.*, *77*, 5399–5406.
- Zhang, T. L., J. G. Luhmann, and C. T. Russell (1990), The solar cycle dependence of the location and shape of the Venus bow shock, *J. Geophys. Res.*, *95*, 14,961–14,967, doi:10.1029/JA095iA09p14961.
- Zhang, T. L., et al. (2006), Magnetic field investigation of the Venus plasma environment: Expected new results from Venus Express, *Planet. Space Sci.*, *54*, 1336–1343.
- Zhang, T. L., et al. (2008), Initial Venus Express magnetic field observations of the Venus bow shock location at solar minimum, *Planet. Space Sci.*, *56*(6), 785–789, doi:10.1016/j.pss.2007.09.012.
- Zhu, D., M. A. Balikhin, M. Gedalin, H. Alleyne, S. A. Billings, Y. Hobara, V. Krasnosel'skikh, M. W. Dunlop, and M. Ruderman (2008), Nonlinear dynamics of foreshock structures: Application of nonlinear autoregressive moving average with exogenous inputs model to Cluster data, *J. Geophys. Res.*, *113*, A04221, doi:10.1029/2007JA012493.